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## OPTICAL SCANNING DEVICE

The present invention relates to an optical scanning device for scanning a first information layer by means of a first radiation beam having a first wavelength and a first polarization, a second information layer by means of a second radiation beam having a second wavelength and a second polarization, and a third information layer by means of a third radiation beam having a third wavelength and a third polarization, wherein said first, second and third wavelengths substantially differ from each other, the device comprising:

a radiation source for emitting said first, second and third radiation beams consecutively or simultaneously,

an objective lens system for converging said first, second and third radiation beams on the positions of said first, second and third information layers, and

a phase structure with a non-periodic stepped profile, arranged in the optical path of said first, second and third radiation beams, the structure including a plurality of steps with different heights for forming said non-periodic stepped profile.

One particular illustrative embodiment of the invention relates to an optical scanning device that is capable of reading data from three different types of optical record carriers, such as compact discs (CDs), conventional digital versatile discs (DVDs) and so-called next generation HD-DVDs.

The present invention also relates to a phase structure for use in an optical scanning device for scanning a first information layer by means of a first radiation beam having a first wavelength and a first polarization, a second information layer by means of a second radiation beam having a second wavelength and a second polarization, and a third information layer by means of a third radiation beam having a third wavelength and a third polarization, wherein said first, second and third wavelengths substantially differ from each other, the structure being arranged in the optical path of said first, second and third radiation beams and having a non-periodic stepped profile.

"Scanning an information layer" refers to scanning by means of a radiation beam for reading information in the information layer ("reading mode"), writing information in the information layer ("writing mode"), and/or erasing information in the information layer ("erase mode"). "Information density" refers to the amount of stored information per unit

area of the information layer. It is determined by, inter alia, the size of the scanning spot formed by the scanning device on the information layer to be scanned. The information density may be increased by decreasing the size of the scanning spot. Since the size of the spot depends, inter alia, on the wavelength  $\lambda$  and the numerical aperture NA of the radiation beam forming the spot, the size of the scanning spot can be decreased by increasing NA and/or by decreasing  $\lambda$ .

In the following a first optical element with an optical axis, e.g. an objective lens, for transforming an object to an image may deteriorate the image by introducing a "wavefront aberration"  $W_{\text{abb}}$ . Wavefront aberrations have different types expressed in the form of the so-called Zernike polynomials with different orders. Wavefront tilt or distortion is an example of a wavefront aberration of the first order. Astigmatism and curvature of field and defocus are two examples of a wavefront aberration of the second order. Coma is an example of a wavefront aberration of the third order. Spherical aberration is an example of a wavefront aberration of the fourth order. It is noted that some wavefront aberrations, such as wavefront tilt, astigmatism and coma, are asymmetric with respect to the optical axis, i.e. dependent on a direction in a plane perpendicular to that axis. Some wavefront modifications, such as defocus and spherical aberration, are symmetric with respect to the optical axis, i.e. independent on any direction in a plane perpendicular to that axis. For more information on the mathematical functions representing the aforementioned wavefront aberrations, see, e.g. the book by M. Born and E. Wolf entitled "Principles of Optics," pp.464-470 (Pergamon Press 6<sup>th</sup> Ed.) (ISBN 0-08-026482-4).

A radiation beam propagating along an optical path has a wavefront  $W$  with a predetermined shape, given by the following equation:

$$\frac{W}{\lambda} = \frac{\Phi}{2\pi} \quad (0a)$$

where " $\lambda$ " and " $\Phi$ " are the wavelength and the phase of the radiation beam, respectively.

In the following a second optical element with an optical axis, e.g. a non-periodic phase structure, may be arranged in the optical path of the radiation beam for introducing a "wavefront modification"  $\Delta W$  in the radiation beam. The wavefront modification  $\Delta W$  is a modification of the shape of the wavefront  $W$ . It may be of a first, second, etc. order of a radius in the cross-section of the radiation beam if the mathematical function describing the wavefront modification  $\Delta W$  has a radial order of three, four, etc., respectively. The wavefront modification  $\Delta W$  may also be "flat"; this means that the second optical element introduces in the radiation beam a constant phase change so that,

after taking modulo  $2\pi$  of the wavefront modification  $\Delta W$ , the resulting wavefront is constant. The term "flat" does not necessarily imply that the wavefront  $W$  exhibits a zero phase change. Furthermore, it derived from Equation (0a) that the wavefront modification  $\Delta W$  may be expressed in the form of a phase change  $\Delta\Phi$  of the radiation beam, given by the following equation:

$$\Delta\Phi = \frac{2\pi}{\lambda} \Delta W \quad (0b)$$

In the following the so-called optical path difference OPD may be calculated for either a wavefront aberration  $W_{abb}$  or a wavefront modification  $\Delta W$ . In the case where the wavefront modification or aberration is symmetric with respect to the optical axis, the root-mean-square value  $OPD_{rms}$  of the optical path difference OPD is given by the following equation:

$$OPD_{rms} = \sqrt{\frac{\int f(r)^2 r dr}{\int r dr} - \left( \frac{\int f(r) r dr}{\int r dr} \right)^2} \quad (0c)$$

where "f" is the mathematical function which describes the wavefront aberration  $W_{abb}$  or the wavefront modification  $\Delta W$  and "r" is the polar coordinate of the polar coordinate system (r,  $\theta$ ) in a plane normal to the optical axis, with the origin of the system is the point of intersection of that plane and the optical axis and extending over the entrance pupil of the corresponding optical element. It is noted that Equation (0c) is applicable to spherical aberration and defocus which are symmetric wavefront aberrations.

In the present description two values  $OPD_{rms,1}$  and  $OPD_{rms,2}$  are "substantially equal" to each other where  $|OPD_{rms,1} - OPD_{rms,2}|$  is less than or equal to, preferably,  $30m\lambda$ , where the value  $30m\lambda$  has been chosen arbitrarily. Also, two values of phase changes  $\Delta\Phi_a$  and  $\Delta\Phi_b$  are "substantially equal" to each other where the respective values  $OPD_{rms,1}$  and  $OPD_{rms,2}$  are "substantially equal" to each other (the relationship between  $\Delta\Phi$  and  $\Delta W$  being given in Equation (0b)). Similarly, two values  $OPD_{rms,1}$  and  $OPD_{rms,2}$  (or two values of phase changes  $\Delta\Phi_a$  and  $\Delta\Phi_b$ ) are "substantially different" from each other where  $|OPD_{rms,1} - OPD_{rms,2}|$  is more than or equal to, preferably,  $30 m\lambda$ , where the value  $30 m\lambda$  has been chosen arbitrarily.

In the following the term "approximate" or "approximation" is used herein, that it is intended to cover a range of possible approximations, the definition including approximations which are in any case sufficient to provide a working embodiment of an

optical scanning device serving the purpose of scanning different types of optical record carriers.

There is currently a need in the field of optical storage for providing optical scanning devices having one optical objective lens for scanning a variety of different optical carriers using different wavelengths of laser radiation, such as a first disc of the so-called BD-format (Blu-ray Disc), a second disc of the so-called DVD-format and a third disc of the so-called CD-format.

For instance, a typical problem is to make an optical scanning device compatible with all currently existing disks, i.e. DVD-format discs and CD-format disc and "HD-DVD"-format discs readout, by means of a first radiation beam with a first wavelength that equals 785nm (to read CD-R), a second radiation beam with a second wavelength that equals 405 nm, and a third radiation beam with a third wavelength that equals 650 nm (to read dual-layer DVD). Due to this plurality of wavelengths, designing a non-periodic phase structure generating predefined wavefronts for each wavelength configuration is difficult. The reason for this is that in designing a non-periodic phase structure (NPS) one makes use of the fact that the phase introduced by a step height  $h$  is different when the wavelength is different. For two wavelength such a structure allows for rather simple designs. It is noted that a method for designing an NPS is known from, e.g., the article by B.H.W. Hendriks, J.E. de Vries and H.P. Urbach, "Application of non-periodic phase structures in optical systems", Appl. Opt. 40 (2001) pp.6548-6560, which describes how to make a objective lens suitable for scanning DVD-format discs and CD-format discs with the aid of an NPS.

It has previously been proposed in, for example, the European Patent application filed on 05.04.2001 with the application number EP 01201255.5, to provide optical scanning devices that are capable of scanning data from HD-DVDs, DVDs and CDs with three radiation beams of different wavelengths, whilst using the same objective lens. Furthermore, it is known in EP 01201255.5 to provide an NPS suitable for three wavelength simultaneously is discussed. The known NPS is a phase structure with a non-periodic stepped profile, arranged in the optical path of the three radiation beams, the structure including a plurality of steps with different heights for forming the non-periodic stepped profile.

Whilst the previously proposed scanning devices provide a solution for situations where three different optical media are illuminated with three associated different wavelengths of light using the same objective lens, they do not provide assistance in providing NPS structures easy to design and manufacture for fixed values of the wavelengths. As a result, the known NPS becomes complex, requiring the making of relatively high steps.

Accordingly, it is an object to an optical scanning device which has a single optical objective lens for scanning a variety of different optical record carriers using at least three radiation beams having three mutually different wavelengths.

This object is reached by an optical scanning device as described in the opening paragraph wherein, according to the invention, said phase structure includes  
5 birefringent material sensitive to said first, second and third polarizations and said stepped profile is designed for introducing a first wavefront modification, a second wavefront modification and a third wavefront modification for said first, second and third wavelengths, respectively, wherein at least one of said first, second and third wavefront modifications is of  
10 a type different from the others and at least one of said first, second and third polarizations differs from the others.

By forming the phase structure from the birefringent material sensitive to the different polarizations of the three radiation beams and by designing the stepped profile for introducing the first wavefront modification, the above-mentioned problem of compatibility  
15 in respect of the first wavelength is then solved. This will be explained in further detail below. Consequently, by comparison with the known NPS, there is for the NPS according to the invention an additional parameter (polarization) which can be used when designing, thereby giving rise to more design freedom. The phase introduced by a step height  $h$  made of a material having refractive index  $n$  at wavelength  $\lambda$  is given by

$$20 \quad \Phi = 2\pi \frac{h(n-1)}{\lambda} \quad (1)$$

Consequently, when the wavelength changes the phase introduced by a step changes. Furthermore, when changing the polarisation and thus changing the refractive index, also a change in phase introduced by the step is generated. Combining both effects for the three wavelengths system, designing NPS's generating predefined wavefronts for each wavelength  
25 is possible with relatively simple stepped structures.

Therefore, an advantage of the optical scanning device provided with the phase structure according to the invention is to scan optical carriers with a plurality of different radiation wavelengths, i.e. to provide a single device for scanning a number of different types of optical record carriers.

30 Another advantage of forming the phase structure according to the invention is to make a phase structure with less amplitude in the height of the steps than in the known phase structure as described in EP 01201255.5.

It is noted that such a phase structure has a non-periodic stepped profile, as opposed to diffraction parts which have each a periodic stepped profile. It is also noted that non-periodic structures and diffraction parts are different from each other in terms of structures and purposes. Thus, an NPS comprises a plurality of steps having different heights so that the NPS has a non-periodic profile. The latter is designed for forming a wavefront modification from a radiation beam incident to the NPS. By contrast, a diffraction part includes a pattern of pattern elements having each one stepped profile. The latter is designed for forming, from a radiation beam incident to the part, a diffracted radiation beam (i.e. a plurality of radiation beams having each a diffraction order "m", i.e. the zeroth order (m=0), the +1<sup>st</sup>-order (m=1), etc., the -1<sup>st</sup>-order (m=-1), etc.) with different transmission efficiencies for different diffraction orders.

In a first embodiment of the optical scanning device according to the invention, said stepped profile is designed for introducing: a second, flat wavefront modification for said second wavelength, and a third, flat wavefront modification for said third wavelength, where at least one of said first, second and third polarisations differs from the others.

In a second embodiment of the optical scanning device according to the invention, said stepped profile is designed for introducing: a second, flat wavefront modification for said second wavelength and, for said third wavelength, a third wavefront modification which substantially is of the same type as said first wavefront modification, where at least one of said first, second and third polarisations differs from the others.

According to another aspect of the invention, the extraordinary refractive index of said birefringent material substantially equals  $1 + \frac{\lambda_c}{\lambda_b}(n_o - 1)$ , where "n<sub>o</sub>" is the ordinary refractive index of said birefringent and "λ<sub>b</sub>" and "λ<sub>c</sub>" are two of said first, second and third wavelengths.

Another object of the invention to provide a phase structure suitable for use in an optical scanning device for scanning a first information layer by means of a first radiation beam having a first wavelength and a first polarization, a second information layer by means of a second radiation beam having a second wavelength and a second polarization, and a third information layer by means of a third radiation beam having a third wavelength and a third polarization, wherein said first, second and third wavelengths substantially differ from each other.

This object is reached by an optical scanning device as described in the opening paragraph wherein, according to the invention, said phase structure includes birefringent material sensitive to said first, second and third polarizations and said stepped profile is designed for introducing a first wavefront modification, a second wavefront  
5 modification and a third wavefront modification for said first, second and third wavelengths, respectively, wherein at least one of said first, second and third wavefront modifications is of a type different from the others and at least one of said first, second and third polarisation differs from the others.

In accordance with another aspect of the invention, there is provided a lens for  
10 use in an optical scanning device for scanning a first information layer by means of a first radiation beam having a first wavelength and a first polarization, a second information layer by means of a second radiation beam having a second wavelength and a second polarization, and a third information layer by means of a third radiation beam having a third wavelength and a third polarization, wherein said first, second and third wavelengths substantially differ  
15 from each other, the lens being provided with a phase structure according to the invention.

The objects, advantages and features of the invention will be apparent from the following, more detailed description of the invention, as illustrated in the accompanying  
20 drawings, in which:

Fig. 1 is a schematic illustration of components of an optical scanning device 1 according to the invention,

Fig. 2 is a schematic illustration of an objective lens for use in the scanning device of Fig. 1,

25 Fig. 3 is a schematic front view of the objective lens of Fig. 2,

Fig. 4 shows a curve representing a wavefront aberration generated by the objective lens shown in Figs. 2 and 3,

Fig. 5 shows a curve representing the step heights of a first embodiment of the NPS shown in Figs. 2 and 3,

30 Fig. 6A shows a curve representing the wavefront modification introduced by the NPS shown in Fig. 5,

Fig. 6B shows a curve representing the combination of the wavefront aberration shown in Fig. 4 and the wavefront modification shown in Fig. 6A, and

Fig. 7 shows a curve representing the step heights of a second embodiment of the NPS shown in Figs. 2 and 3.

5 Fig. 1 is a schematic illustration of the optical components of an optical scanning device 1 according to one embodiment of the invention, for scanning a first information layer 2'' of a first optical record carrier 3'' by means of a first radiation beam 4''.

By way of illustration, the optical record carrier 3'' includes a transparent layer 5'' on one side of which the information layer 2'' is arranged. The side of the information  
10 layer facing away from the transparent layer 5'' is protected from environmental influences by a protective layer 6''. The transparent layer 5'' acts as a substrate for the optical record carrier 3'' by providing mechanical support for the information layer 2''. Alternatively, the transparent layer 5'' may have the sole function of protecting the information layer 2'', while the mechanical support is provided by a layer on the other side of the information layer 2'',  
15 for instance by the protective layer 6'' or by an additional information layer and transparent layer connected to the uppermost information layer. It is noted that the information layer has a first information layer depth 27'' that corresponds to, in this embodiment as shown in Fig. 1, to the thickness of the transparent layer 5''. The information layer 2'' is a surface of the carrier 3''. That surface contains at least one track, i.e. a path to be followed by the spot of a  
20 focused radiation on which path optically-readable marks are arranged to represent information. The marks may be, e.g., in the form of pits or areas with a reflection coefficient or a direction of magnetization different from the surroundings. In the case where the optical record carrier 3'' has the shape of a disc, the following is defined with respect to a given track: the "radial direction" is the direction of a reference axis, the X-axis, between the track  
25 and the center of the disc and the "tangential direction" is the direction of another axis, the Y-axis, that is tangential to the track and perpendicular to the X-axis.

As shown in Fig.1, the optical scanning device 1 includes a radiation source 7, a collimator lens 18, a beam splitter 9, an objective lens system 8 having an optical axis 19, a phase structure or non-periodic structure (NPS) 24, and a detection system 10. Furthermore,  
30 the optical scanning device 1 includes a servocircuit 11, a focus actuator 12, a radial actuator 13, and an information processing unit 14 for error correction.

In the following "Z-axis" corresponds to the optical axis 19 of the objective lens system 8. It is noted that (X, Y, Z) is an orthogonal base.



The radiation source 7 is arranged for consecutively or simultaneously supplying the radiation beam 4'' and two other radiation beams 4 and 4' (not shown in Fig. 1). For example, the radiation source 7 may comprise either a tunable semiconductor laser for consecutively supplying the radiation beams 4'', 4 and 4' or three semiconductor lasers for simultaneously supplying these radiation beams. Furthermore, the radiation beam 4'' has a first wavelength  $\lambda_3$  and a first polarization  $p_3$ , the radiation beam 4 has a second wavelength  $\lambda_1$  and a second polarization  $p_1$ , and the radiation beam 4' has a third wavelength  $\lambda_2$  and a third polarization  $p_2$ . Examples of the wavelengths  $\lambda_1$ ,  $\lambda_2$  and  $\lambda_3$  and the polarizations  $p_1$ ,  $p_2$  and  $p_3$  will be given where the wavelengths  $\lambda_1$ ,  $\lambda_2$  and  $\lambda_3$  substantially differ from each other and the polarization  $p_3$  differs from at least one of the polarizations  $p_1$  and  $p_2$ . It is noted in the present description that two wavelengths  $\lambda_a$  and  $\lambda_b$  are substantially different from each other where  $|\lambda_a - \lambda_b|$  is equal to or higher than, preferably, 10nm and, more preferably, 20nm, where the values 10 and 20nm are a matter of a purely arbitrary choice.

The collimator lens 18 is arranged on the optical axis 19 for transforming the radiation beam 4'' into a first substantially collimated beam 20''. Similarly, it transforms the radiation beams 4 and 4' into a second substantially collimated beam 20 and a third substantially collimated beam 20' (not shown in Fig. 1).

The beam splitter 9 is arranged for transmitting the collimated radiation beams 20'', 20 and 20' toward the objective lens system 8. Preferably, the beam splitter 9 is formed with a plane parallel plate that is tilted with an angle  $\alpha$  with respect to the Z-axis and, more preferably,  $\alpha=45^\circ$ .

The objective lens system 8 is arranged for transforming the collimated radiation beam 20'' to a first focused radiation beam 15'' so as to form a first scanning spot 16'' in the position of the information layer 2''. Similarly, the objective lens system 8 transforms the collimated radiation beams 20 and 20' as explained below.

In this embodiment, the objective lens system 8 includes an objective lens 17 provided with the NPS 24.

The NPS 24 includes birefringent material having an extraordinary refractive index  $n_e$  and an ordinary refractive index  $n_o$ . In the following the change in refractive index due to difference in wavelength is neglected and therefore the refractive indices  $n_e$  and  $n_o$  are approximately independent of the wavelength. In this embodiment, and by way of illustration only, the birefringent material is C6M/E7 50/50 (in % by weight) with  $n_o=1.51$  and  $n_e=1.70$ .

Alternatively, for example, the birefringent material may be C6M/C3M/E7 40/10/50 (in % by weight) with  $n_o=1.55$  and  $n_e=1.69$ . The codes used refer to the following substances:

E7: 51% C5H11cyanobiphenyl, 25% C5H15cyanobiphenyl, 16% C8H17cyanobiphenyl, 8% C5H11 cyanotriphenyl;

5 C3M: 4-(6-acryloyloxypropyloxy)benzoyloxy-2-methylphenyl 4-(6-acryloyloxypropyloxy)benzoate;

C6M: 4-(6-acryloyloxyhexyloxy)benzoyloxy-2-methylphenyl 4-(6-acryloyloxyhexyloxy)benzoate.

The NPS 24 is aligned such that the optic axis of the birefringent material is  
 10 along the Z-axis. It is also aligned such that its refractive index equals  $n_e$  when traversed by a radiation beam having a polarisation along the X-axis and  $n_o$  when traversed by a radiation beam having a polarisation along the Y-axis. In the following the polarization of a radiation beam is called " $p_e$ " and " $p_o$ " where aligned with the X-axis and the Y-axis, respectively. Thus, where the polarization  $p_1$ ,  $p_2$  or  $p_3$  equals  $p_e$ , the refractive index of the birefringent  
 15 material equals  $n_e$  and, where the polarization  $p_1$ ,  $p_2$  or  $p_3$  equals  $p_o$ , the refractive index of the birefringent material equals  $n_o$ . In other words, the birefringent NPS 24 so aligned is sensitive to the polarizations  $p_1$ ,  $p_2$  and  $p_3$ . The NPS 24 will be described in further detail.

During scanning, the record carrier 3" rotates on a spindle (not shown in Fig. 1) and the information layer 2" is then scanned through the transparent layer 5". The  
 20 focused radiation beam 15" reflects on the information layer 2", thereby forming a reflected beam 21" which returns on the optical path of the forward converging beam 15". The objective lens system 8 transforms the reflected radiation beam 21" to a reflected collimated radiation beam 22". The beam splitter 9 separates the forward radiation beam 20" from the reflected radiation beam 22" by transmitting at least a part of the reflected radiation beam 22"  
 25 towards the detection system 10.

The detection system 6 includes a convergent lens 25 and a quadrant detector 23 which are arranged for capturing said part of the reflected radiation beam 22" and converting it to one or more electrical signals. One of the signals is an information signal  $I_{data}$ , the value of which represents the information scanned on the information layer 2". The  
 30 information signal  $I_{data}$  is processed by the information processing unit 14 for error correction. Other signals from the detection system 10 are a focus error signal  $I_{focus}$  and a radial tracking error signal  $I_{radial}$ . The signal  $I_{focus}$  represents the axial difference in height along the Z-axis between the scanning spot 16" and the position of the information layer 2". Preferably, this signal is formed by the "astigmatic method" which is known from, inter alia, the book by G.

Bouwhuis, J. Braat, A. Huijser et al, entitled "Principles of Optical Disc Systems," pp.75-80 (Adam Hilger 1985) (ISBN 0-85274-785-3). The radial tracking error signal  $I_{\text{radial}}$  represents the distance in the XY-plane of the information layer 2" between the scanning spot 16" and the center of a track in the information layer 2" to be followed by the scanning spot 16".

5 Preferably, this signal is formed from the "radial push-pull method" which is known from, inter alia, the book by G. Bouwhuis, pp.70-73.

The servocircuit 11 is arranged for, in response to the signals  $I_{\text{focus}}$  and  $I_{\text{radial}}$ , providing servo control signals  $I_{\text{control}}$  for controlling the focus actuator 12 and the radial actuator 13, respectively. The focus actuator 12 controls the position of the objective lens 17  
10 along the Z-axis, thereby controlling the position of the scanning spot 16" such that it coincides substantially with the plane of the information layer 2". The radial actuator 13 controls the position of the objective lens 17 along the X-axis, thereby controlling the radial position of the scanning spot 16" such that it coincides substantially with the center line of the track to be followed in the information layer 2".

15 Fig. 2 is a schematic illustration of the objective lens 17 for use in the scanning device 1 described above.

The objective lens 17 is arranged for transforming the collimated radiation beam 20" to the focused radiation beam 15", having a first numerical aperture  $NA_3$ , so as to form the scanning spot 16". In other words, the optical scanning device 1 is capable of  
20 scanning the first information layer 2" by means of the radiation beam 15" having the wavelength  $\lambda_3$ , the polarization  $p_3$  and the numerical aperture  $NA_3$ .

Furthermore, the optical scanning device 1 is also capable of scanning a second information layer 2 of a second optical record carrier 3 by means of the radiation beam 4 and a third information layer 2' of a third optical record carrier 3' by means of the  
25 radiation beam 4'. Thus, the objective lens 17 transforms the collimated radiation beam 20 to a second focused radiation beam 15, having a second numerical aperture  $NA_1$ , so as to form a second scanning spot 16 in the position of the information layer 2. The objective lens 17 also transforms the collimated radiation beam 20' to a third focused radiation beam 15', having a third numerical aperture  $NA_2$ , so as to form a third scanning spot 16' in the position of the  
30 information layer 2'.

Similarly to the optical record carrier 3", the optical record carrier 3 includes a second transparent layer 5 on one side of which the information layer 2 is arranged with a second information layer depth 27, and the optical record carrier 3' includes a third

transparent layer 5' on one side of which the information layer 2' is arranged with a third information layer depth 27'.

It is noted that scanning information layers of the record carriers 3, 3' and 3'' of different formats is achieved by forming the objective lens 17 as a hybrid lens, i.e. a lens combining an NPS and refractive elements, used in an infinite-conjugate mode. Such a hybrid lens can be formed by applying a stepped profile on the entrance surface of the lens 17, for example by a lithographic process using the photopolymerisation of, e.g., an UV curing lacquer, thereby advantageously resulting in the NPS 24 to be easy to make. Alternatively, such a hybrid lens can be made by diamond turning.

In the embodiment shown in Figs. 1 and 2, the objective lens 17 is formed as a convex-convex lens; however, other lens element types such as plano-convex or convex-concave lenses can be used. In this embodiment, the NPS 24 is arranged on the side of a first objective lens 17 facing the radiation source 7 (referred to herein as the "entrance face").

Alternatively, the NPS 24 is arranged on the other surface of the lens 17 (referred to herein as the "exit face"). Also alternatively, the objective lens 17 is, for example, a refractive objective lens element provided with a planar lens element forming the NPS 24. Also alternatively, the NPS 24 is provided on an optical element separate from the objective lens system 8, for example on a beam splitter or a quarter wavelength plate.

Also alternatively, whilst the objective lens 17 is in this embodiment a single lens, it may be a compound lens containing two or more lens element.

Fig. 3 is a schematic view of the entrance surface (also called "front view") of the objective lens 17 shown in Fig. 2, illustrating the NPS 24.

The NPS 24 includes a plurality of steps "j" with different heights "h<sub>j</sub>" for forming the non-periodic stepped profile. In the following "h" is the step height of the stepped profile, which is a function dependent on x. In the case of the stepped-profile approximation, the step height h is given by the following function:

$$h(x) = h_j \quad \text{for } j-1 \leq x \leq j \quad (2a)$$

where "h<sub>j</sub>" is the step height of the step j, which is a constant parameter. In the following "zone" is the length of a step along the X-axis.

The stepped profile is designed, i.e. the step height h<sub>j</sub> are chosen, for introducing a first wavefront modification ΔW<sub>3</sub> (and therefore a first phase change ΔΦ<sub>3</sub>) at the wavelength λ<sub>3</sub>, a second wavefront modification ΔW<sub>1</sub> (and therefore a second phase change ΔΦ<sub>1</sub>) at the wavelength λ<sub>1</sub>, and a third wavefront modification ΔW<sub>2</sub> (and therefore a third phase change ΔΦ<sub>2</sub>) at the wavelength λ<sub>2</sub>. In other words, the stepped profile is designed

so as to introduce the wavefront modifications  $\Delta W_1, \Delta W_2$  and  $\Delta W_3$  in the radiation beams 15, 15' and 15'' where these wavefront modifications are either flat of a type of a symmetric aberration.

In the following and by way of illustration only the wavefront modification  $\Delta W_1$  is flat. Thus, the step heights  $h_j$  are chosen so that the phase change  $\Delta\Phi_1$  substantially equals a multiple of  $2\pi$ , i.e. substantially equal zero modulo  $2\pi$ . In this embodiment the wavelength  $\lambda_1$  is said to be the design wavelength  $\lambda_{ref}$ . In other words,

$$\lambda_{ref} = \lambda_1 \quad (2b)$$

$$\Delta\Phi_1 \equiv 0 \pmod{2\pi}. \quad (2c)$$

This is achieved when each step height  $h_j$  is a multiple of a reference height  $h_{ref}$  which is dependent on the design wavelength  $\lambda_{ref}$  (i.e. the wavelength  $\lambda_1$ ) as follows:

$$h_{ref} = \frac{\lambda_{ref}}{n - n_0} \quad (3)$$

where "n" is the refractive index of the NPS 24 and  $n_0$  is the refractive index of the adjacent medium that is, in the following and by way of illustration only, air, i.e.  $n_0=1$ .

It is noted that the reference height  $h_{ref}$  is substantially constant, in the case where the NPS 24 is provided on a plane surface (e.g. on a plane parallel plate). Furthermore, in the case when the NPS 24 is provided on a curved surface (e.g. that of a lens), the NPS 24 may be adjusted over the length of the step so as to generate phase changes that are substantially equally to multiple of  $2\pi$ .

Since the NPS 24 is made of birefringent material, its refractive index  $n$  equals  $n_e$  when the polarization of the radiation beam traversing the NPS 24 equals  $p_e$  and equals  $n_o$  when the polarization of the radiation beam traversing the NPS 24 equals  $p_o$ . Consequently, the reference height  $h_{ref}$  is dependent on the reference wavelength  $\lambda_{ref}$  and also the polarization  $p_{ref}$  of the reference wavelength  $\lambda_{ref}$  and in the following it is also referred to as " $h_{ref}(\lambda_{ref}, p_{ref})$ ". Similarly, the phase changes  $\Delta\Phi_1, \Delta\Phi_2$  and  $\Delta\Phi_3$  are also dependent the respective polarizations  $p_1, p_2$  and  $p_3$  and in the following they are also referred to as " $\Delta\Phi_1(p_1)$ ", " $\Delta\Phi_2(p_2)$ " and " $\Delta\Phi_3(p_3)$ ".

Consequently, it follows from Equations (2b) and (3) that:

$$h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_e) = \frac{\lambda_1}{n_e - n_0} \quad (4a)$$

$$h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_o) = \frac{\lambda_1}{n_o - n_0} \quad (4b)$$

Accordingly, in the case where, e.g.,  $n_o=1.50$ ,  $n_e=1.62$  and  $\lambda_1=405\text{nm}$ , the following is obtained from Equations (4a) and (4b):

$$h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_e)=0.653\mu\text{m} \text{ and}$$

$$h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_o)=0.810\mu\text{m}.$$

- 5 It is also noted that, while a step height  $h_j$  introduces the value  $\Delta\Phi_1(p_1)$  (substantially equal to zero modulo  $2\pi$ ) for the radiation beam 15, it introduces the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  for the radiation beams 15' and 15'', respectively, as follows:

$$\Delta\Phi_2(p_2=p_e) = 2\pi \frac{n_e - n_o}{\lambda_2} h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1) \quad (5a)$$

$$\Delta\Phi_2(p_2=p_o) = 2\pi \frac{n_o - n_o}{\lambda_2} h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1) \quad (5b)$$

$$10 \quad \Delta\Phi_3(p_3=p_e) = 2\pi \frac{n_e - n_o}{\lambda_3} h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1) \quad (5c)$$

$$\Delta\Phi_3(p_3=p_o) = 2\pi \frac{n_o - n_o}{\lambda_3} h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1) \quad (5d)$$

- Table I shows the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  where the radiation beams 15' and 15'' traverse the step height  $h_j$  which equals either  $h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_e)$  or  $h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_o)$ , in the cases where the polarizations  $p_2$  and  $p_3$  equal  $p_e$  and/or  $p_o$ . The values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  have been calculated from Equations (4a), (4b) and (5a) to (5d) with, e.g.,  $n_o=1.50$ ,  $n_e=1.62$ ,  $\lambda_1=405\text{nm}$ ,  $\lambda_2=650\text{nm}$  and  $\lambda_3=785\text{nm}$ .

Table I

		$\Delta\Phi_2(p_2)/2\pi$ (modulo 1)		$\Delta\Phi_3(p_3)/2\pi$ (modulo 1)	
		$p_2=p_e$	$p_2=p_o$	$p_3=p_e$	$p_3=p_o$
$h_j=h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1)$	$p_1=p_e$	0.623	0.502	0.516	0.416
	$p_1=p_o$	0.773	0.623	0.640	0.516

- It is further noted that a step height  $h_j$  equal to a multiple of  $h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1)$  introduces the value  $\Delta\Phi_1(p_1)$  that equals zero modulo  $2\pi$  for the radiation beam 15 and the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  that each equal one among a limited number of possible values. In the following “ $\#\Delta\Phi_2$ ” and “ $\#\Delta\Phi_3$ ” are such limited numbers for the values of the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$ , respectively. Similarly to the phase changes  $\Delta\Phi_1$ ,  $\Delta\Phi_2$  and  $\Delta\Phi_3$ ,

the limited numbers  $\#\Delta\Phi_2$  and  $\#\Delta\Phi_3$  are also dependent the respective polarizations  $p_2$  and  $p_3$  and in the following they are also referred to as “ $\#\Delta\Phi_2(p_2)$ ” and “ $\#\Delta\Phi_3(p_3)$ ”, The limited numbers  $\#\Delta\Phi_2(p_2)$  and  $\#\Delta\Phi_3(p_3)$  have been calculated based on the theory of Continued Fractions, as known from, e.g., the European patent application filed on 05.04.2001 under the application number 01201255.5.

By way of illustration only, the calculation of the limited numbers  $\#\Delta\Phi_3(p_3)$  is now described in a first case where the polarizations  $p_1$  and  $p_3$  are identical, e.g.  $p_1=p_0$  and  $p_3=p_0$ , and a second case where the polarization  $p_1$  differs from the polarization  $p_3$ , e.g.  $p_1=p_0$  and  $p_3=p_e$ . With reference to said European patent application filed under the application number 01201255.5, the following is defined:

$$a_0 = \frac{H_i}{H_l} \quad (6a)$$

$$b_0 = \text{Int}[a_0] \quad (6b)$$

$$a_1 = a_0 - b_0 \quad (6c)$$

$$b_m = \text{Int}\left[\frac{1}{a_m}\right] \quad (6d)$$

$$a_{m+1} = \frac{1}{a_m} - b_m \quad (6e)$$

$$CF_m = \{b_0, b_1 \dots b_m\} \quad (6f)$$

where  $H_l = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_1)$ ,  $H_i = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_3, p_{\text{ref}}=p_3)$  and “m” is an integer equal to or higher than 1.

In the first case where  $p_1=p_0$  and  $p_3=p_0$  and where, e.g.,  $n_o=1.5$ ,  $n_e=1.62$ ,  $\lambda_1=405\text{nm}$  and  $\lambda_3=785\text{nm}$ , the following is obtained from Equations (6a) to (6e):

$$H_l = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_0) = \frac{\lambda_1}{n_o - n_e} = 0.810\mu\text{m}$$

$$H_i = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_3, p_{\text{ref}}=p_0) = \frac{\lambda_3}{n_o - n_e} = 1.570\mu\text{m}$$

$$a_0 = 0.516$$

$$b_0 = 0$$

$$a_1 = 0.516$$

$$b_1 = 1$$

$$a_2 = 0.938$$

$$b_2 = 1$$

$$CF_2 = 0 + \frac{1}{1 + \frac{1}{1}} = \frac{1}{2}$$

Thus,  $CF_2$  substantially equals  $a_0$ , i.e. the following is met:  $|CF_2 - a_0| = 0.016 < 0.02$  where 0.02 is a value chosen purely arbitrarily. As a result, it is found that the limited number  $\# \Delta \Phi_3(p_3=p_0)$  is equal to 2 where  $p_1=p_0$ .

In the second case where  $p_1=p_0$  and  $p_3=p_e$ , and where, e.g.,  $n_o=1.50$ ,  $n_e=1.62$ ,

5  $\lambda_1=405\text{nm}$  and  $\lambda_3=785\text{nm}$ , the following is obtained from Equations (6a) to (6e):

$$H_1 = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_1, p_{\text{ref}}=p_0) = \frac{\lambda_1}{n_o - n_0} = 0.810\mu\text{m}$$

$$H_i = h_{\text{ref}}(\lambda_{\text{ref}}=\lambda_3, p_{\text{ref}}=p_e) = \frac{\lambda_3}{n_e - n_0} = 1.266\mu\text{m}$$

$$a_0=0.640$$

$$b_0=0$$

$$10 \quad a_1=0.640$$

$$b_1=1$$

$$a_2=0.563$$

$$b_2=1$$

$$a_3=0.776$$

$$15 \quad b_3=1$$

$$a_4=0.288$$

$$b_4=3$$

$$CF_4 = 0 + \frac{1}{1 + \frac{1}{1 + \frac{1}{1 + \frac{1}{3}}}} = \frac{7}{11}$$

Thus,  $CF_4$  substantially equals  $a_0$ , i.e. the following is met:  $|CF_4 - a_0| = 0.004 < 0.02$ . As a

20 result, it is found that the limited number  $\# \Delta \Phi_3(p_3=p_e)$  is equal to 11 where  $p_1=p_0$ .

Table II shows the limited numbers  $\# \Delta \Phi(\lambda=\lambda_2, p=p_2)$  and  $\# \Delta \Phi(\lambda=\lambda_3, p=p_3)$  in respect of a step height  $h_j$  equal to  $h_{\text{ref}}(\lambda=\lambda_1, p=p_e)$  and  $h_{\text{ref}}(\lambda=\lambda_1, p=p_0)$  and in the cases where the polarizations  $p_2$  and  $p_3$  equal  $p_e$  and/or  $p_0$ . These limited numbers have been calculated on the theory of Continued Fractions as described above.



Table II:

		# $\Delta\Phi_2(p_2)$		# $\Delta\Phi_3(p_3)$	
		$p_2=p_e$	$p_2=p_o$	$p_3=p_e$	$p_3=p_o$
$h_j=h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$	$p_1=p_e$	8	2	2	5
	$p_1=p_o$	9	8	11	2

It is noted in Tables I and II that if the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are identical, one of the limited numbers  $\# \Delta\Phi_2(p_2)$  and  $\# \Delta\Phi_3(p_3)$  equals 2, i.e. only two different values (zero and  $\pi$  modulo  $2\pi$ ) can be chosen for the corresponding phase changes. This does not allow a substantial degree of freedom for designing the NPS 24 in respect of the corresponding radiation beam.

By contrast, it is also noted in Tables I and II that if at least one of the polarizations  $p_1$ ,  $p_2$ ,  $p_3$  differs from the others, at least three different values can be chosen for  $\Delta\Phi_2(p_2)$  and/or  $\Delta\Phi_3(p_3)$ . The possibility for choosing the phase changes from at least three possible values allows to make an efficient NPS for each of the radiation beams 15, 15' and 15". Furthermore, this advantageously allows to design the stepped profile with a relatively low number of steps, typically less than 40 steps, since a stepped profile with a high number of steps (typically, 50 or more steps) is of less practical use.

Two embodiments of the stepped profile are now described where the wavefront modification  $\Delta W_3$  is of the type of a symmetric aberration and the wavefront modification  $\Delta W_2$  is flat in the first embodiment and of the type of a symmetric aberration in the second embodiment.

In the first embodiment and by way of illustration only the optical record carriers 3, 3' and 3" are a "HD-DVD"-format disc, a DVD-format disc and a CD-format disc, respectively. Firstly, the wavelength  $\lambda_1$  is comprised in the range between 365 and 445nm and, preferably, 405nm. The wavelength  $\lambda_2$  is comprised in the range between 620 and 700nm and, preferably, 650nm. The wavelength  $\lambda_3$  is comprised in the range between 740 and 820nm and, preferably, 785nm. Secondly, the numerical aperture  $NA_1$  equals about 0.6 in the reading mode and is above 0.6, preferably 0.65, in the writing mode. The numerical aperture  $NA_2$  equals about 0.6 in the reading mode and is above 0.6, preferably 0.65, in the

writing mode. The numerical aperture  $NA_3$  is below 0.5, preferably 0.45. Thirdly, the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are as follows:  $p_1=p_e$ ,  $p_2=p_o$  and  $p_3=p_o$ .

In the first embodiment, the objective lens 17 is a plano-aspherical element (as shown in Fig. 2). The objective lens 17 has a thickness of 2.412mm along on the Z-axis (i.e. the direction of its optical axis) and an entrance pupil with a diameter of 3.3mm. The numerical aperture of the objective lens 17 is equal to 0.6 at the wavelength  $\lambda_1$  (=405nm), to 0.6 at the wavelength  $\lambda_2$  (=650nm), and to 0.45 at the wavelength  $\lambda_3$  (=785nm). The lens body of the objective lens is made of LAFN28 Schott glass with a refractive index that is equal to 1.7998 at the wavelength  $\lambda_1$  (=405nm), to 1.7688 at the wavelength  $\lambda_2$  (=650nm), and to 1.7625 at the wavelength  $\lambda_3$  (=785nm). The convex surface of the lens body which is directed towards the collimator lens 18 has a radius of 2.28mm. The surface of the objective lens 17 facing the record carrier is flat. The aspherical shape is realized in a thin layer of acryl on top of the glass body. The lacquer has a refractive index equal to 1.5945 at the wavelength  $\lambda_1$  (=405nm), to 1.5646 at the wavelength  $\lambda_2$  (=650nm), and to 1.5588 at the wavelength  $\lambda_3$  (=785nm). The thickness of this layer on the optical axis is 17 $\mu$ m. The rotational symmetric aspherical shape is defined by a function  $H(r)$  as follows:

$$H(r) = \sum_{i=1}^5 B_{2i} r^{2i} \quad (7)$$

where “ $H(r)$ ” is the position of the surface along the optical axis of the lens 17 in millimeters, “ $r$ ” is the distance to the optical axis in millimeters, and “ $B_k$ ” are the coefficient of the  $k$ -th power of  $H(r)$ . The value of the coefficients  $B_2$  until  $B_{10}$  are 0.238864, 0.0050434889, 7.3344175  $10^{-5}$ , -7.0483109  $10^{-5}$ , -4.7795094  $10^{-6}$ , respectively. The free working distance, i.e. the distance between the objective lens 17 and the optical record carrier, is equal: to 0.9676mm at the wavelength  $\lambda_1$  (=405nm) for a DHD-DVD-format disk having a cover layer thickness of 0.6mm, to 1.044mm at the wavelength  $\lambda_2$  (=650nm) for a DVD-format disk having a cover layer thickness of 0.6mm, and to 0.6917mm at the wavelength  $\lambda_3$  (=785nm) for a CD-format disk having a cover layer thickness of 1.2mm. The cover layer thickness of the disk is made of polycarbonate with refractive index equal to 1.6188 at the wavelength  $\lambda_1$  (=405nm), to 1.5806 at the wavelength  $\lambda_2$  (=650nm) and to 1.5731 at the wavelength  $\lambda_3$  (=785nm). The objective lens 17 is designed in such a way that, when scanning a HD-DVD-format disk at the wavelength  $\lambda_1$  (=405nm) and a DVD-format disc at the wavelength  $\lambda_2$  (=650nm), no spherochromatism is introduced. It is noted that the objective lens 17 is compatible with the “HD-DVD”-format and the DVD-format. In order to make the objective

lens suitable for scanning a CD-format disk, the amount of spherical aberration  $W_{abb}$  arising due to the difference of cover layer thickness and spherochromatism has to be compensated. Spherical aberration can be expressed in the form of the Zernike polynomials. For further information, see e.g. M. Born and E. Wolf, "Principles of Optics," p.469-470 (6<sup>th</sup> ed.)

(Pergamon Press) (ISBN 0-08-09482-4). It is noted that, knowing the shape of the objective lens 17 from Equation (7), the amount of spherical aberration  $W_{abb}$  can be determined by ray-tracing simulations. Fig. 4 shows a curve 81 representing the wavefront aberration  $W_{abb}$  generated by the objective lens 17 according to Equation (7). It is noted in Fig. 4 that " $r_o$ " is the pupil radius of the face of the objective lens 17, which is provided with the NPS 24.

Therefore, in the first embodiment, the stepped profile is designed for compensating the wavefront aberration  $W_{abb}$  at the wavelength  $\lambda_3$ . Thus the step heights  $h_j$  are to be chosen such that the wavefront modifications  $\Delta W_1$  and  $\Delta W_2$  are substantially flat and such that the wavefront modification meets the following:

$$\Delta W_3 \approx -W_{abb} \quad (8)$$

It is noted that the wavefront modifications  $\Delta W_1$  and  $\Delta W_2$  may substantially differ from each other by a substantially constant phase difference.

Accordingly, the step heights  $h_j$  are chosen such that both the phase changes  $\Delta\Phi_1(p_1)$  and  $\Delta\Phi_2(p_2)$  are substantially equal to a constant (e.g. zero) modulo  $2\pi$ , where the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_1(p_1)$  may substantially differ from each other, and such that the sum of the wavefront modification  $\Delta W_3$  and the wavefront aberration  $W_{abb}$  substantially equals zero. By way of illustration only, an example of the first embodiment of the stepped profile is described in the following where the stepped profile includes five steps.

Firstly, Table III shows the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  introduced by step heights that equal  $qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$  where  $p_1=p_e$  and " $q$ " is an integer. These values are found from Table I where the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  are known a step height that equals  $h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$  where  $p_1=p_e$ , i.e. for  $q=1$ .

Table III:

q	$\Delta\Phi_2(p_2)/2\pi$ (modulo 1) $p_2=p_0$	$\Delta\Phi_3(p_3)/2\pi$ (modulo 1) $p_3=p_0$
1	0.502	0.416
2	0.004	0.832
3	0.506	0.248
4	0.008	0.664
5	0.510	0.080
6	0.012	0.496
7	0.514	0.912
8	0.016	0.328
9	0.518	0.744
10	0.020	0.160
11	0.522	0.576
12	0.026	0.992

It is noted in Table III that the phase change  $\Delta\Phi_2(p_2)$  is substantially equal to zero or  $\pi$  modulo  $2\pi$  and that the phase change  $\Delta\Phi_3(p_3)$  has substantially 5 different values modulo  $2\pi$ . This is consistent with Table II where  $\#\Delta\Phi_2(p_2)=2$  for  $p_2=p_0$  and  $\#\Delta\Phi_3(p_3)=5$  for  $p_3=p_0$ .

It is also noted that, since the polarization  $p_3$  differs from the polarization  $p_1$ , at least three different values of the phases changes  $\Delta\Phi_3(p_3)$  can be chosen, thereby resulting in allowing the design of the stepped profile with a relatively low number of steps, typically less than 40 steps, since a stepped profile with a high number of steps (typically, 50 or more steps) is of less practical use.

Secondly, Table IV shows the "optimized zones" of the step height  $h_j$  ( $=qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$ ) where  $p_1=p_e$  and the values of the phase change  $\Delta\Phi_3(p_3)/2\pi$ , which are determined from Table III with  $p_3=p_0$  and the wavefront aberration  $W_{abb}$  (see Fig. 4) according to the method known from said article by B.H.W. Hendriks et al.. Table IV also shows, for the step height  $h_j$ , the values of the phase change  $\Delta\Phi_2(p_2)$  for approximating the flat wavefront modification  $\Delta W_2$  according to Table III where  $p_2=p_0$ .

Table IV:

	Zones (mm)	q	$h_j(\mu\text{m})$	$\Delta\Phi_2(p_2)$ (mod. $2\pi$ ) $p_2=p_0$	$\Delta\Phi_3(p_3)$ (mod. $2\pi$ ) $p_3=p_0$
j=1	0.00-0.40	0	0.000	0.0000	0.000
j=2	0.40-0.59	10	6.530	0.1256	1.005
j=3	0.59-1.10	8	5.224	0.1005	2.061
j=4	1.10-1.20	10	6.530	0.1256	1.005
j=5	1.20-1.26	0	0.000	0.0000	0.000

It is also noted in Table IV that, due to the possibility to choose the value of the refractive index based on the polarizations of the radiation beams, the NPS has an advantageous stepped profile with a difference in the step heights of only 6.53 $\mu\text{m}$ . By contrast, the NPS known from said patent application EP 01201255.5 has a difference in the step heights of more than 16 $\mu\text{m}$ , thereby resulting in the known NPS which is difficult to make.

Fig. 5 shows a curve 80 representing the step height  $h(x)$  of the NPS 24 according to Table IV. It is noted in respect of the curve 80 that the stepped profile is designed such that the relative step heights  $h_{j+1}-h_j$  between adjacent steps include a relative step height having an optical path substantially equal to  $a\lambda_l$ , wherein "a" is an integer and  $a>1$  and " $\lambda_l$ " is the design wavelength. In other words, such a relative step height is higher than the reference height  $h_{\text{ref}}(\lambda=\lambda_l, p=p_1)$ .

Fig. 6A shows a curve 82 representing the wavefront modification  $\Delta W_3$  introduced by the NPS shown in Fig. 5 for compensating the wavefront aberration  $W_{\text{abb}}$ . It is noted in Fig. 6A that the reference "j" corresponds to the steps as defined in relation with Fig. 5.

By comparison, Fig. 6B shows a curve 83 representing the combination of the wavefront aberration shown in Fig. 4 and the wavefront modification shown in Fig. 6A.

By referring again to Table IV, it is also noted that the phase changes  $\Delta\Phi_2(p_2)$  are substantially equal to zero, thereby introducing the flat wavefront modification  $\Delta W_2$ , and that the phase change  $\Delta\Phi_3(p_3)$  associated with the corresponding optimized zones approximates the wavefront aberration  $W_{\text{abb}}$  (here, spherical aberration).

Table V shows the values  $\text{OPD}_{\text{rms}}[W_{\text{abb}}+\Delta W_i]$  for the wavefront modifications  $\Delta W_1$ ,  $\Delta W_2$  and  $\Delta W_3$  where the radiation beams 15, 15' and 15'' (at the respective

wavelengths and polarizations) traverse the NPS for compensating the wavefront aberration  $W_{abb}$  according to Table IV (and shown in Fig. 4). Table V also shows the values

$OPD_{rms}[W_{abb}]$  associated with the wavefront aberration  $W_{abb}$  (i.e. without the correction of the NPS 24 according to Table IV). The values  $OPD_{rms}[W_{abb}+\Delta W_i]$  and  $OPD_{rms}[W_{abb}]$  have

5 been calculated from ray-tracing simulations.

Table V:

	$OPD_{rms}[W_{abb}+\Delta W_i]$	$OPD_{rms}[W_{abb}]$
i=1 ( $p_1=p_e$ )	17.9m $\lambda$	17.9m $\lambda$
i=2 ( $p_2=p_o$ )	8.6m $\lambda$	3.2m $\lambda$
i=3 ( $p_3=p_o$ )	43.8m $\lambda$	134.1m $\lambda$

10 It is noted in Table V that the three values  $OPD_{rms}[W_{abb}+\Delta W_i]$  are below the diffraction limit, i.e. less than 70 m $\lambda$ , for the NPS 24 according to Table IV, thereby allowing any format of optical record carriers to be scanned.

As an alternative of the first embodiment of the stepped profile, the values of the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_1(p_1)$  are substantially equal to each other, where the polarization  $p_1$  different from the polarization  $p_2$ , i.e.:

$$15 \quad \Delta\Phi_2(p_2) = \Delta\Phi_1(p_1) \quad (9)$$

In the case where  $p_1=p_o$ ,  $p_2=p_e$  and  $p_3=p_e$  it derives from Equations (0c), (5b), (5c) and (9) that:

$$\frac{\lambda_2}{n_e - 1} = \frac{\lambda_1}{n_o - 1} \quad (10)$$

It follows from Equation (10) that:

$$20 \quad n_e = 1 + \frac{\lambda_2}{\lambda_1} (n_o - 1) \quad (11)$$

Thus, for example, in the case where  $n_o=1.5$ ,  $\lambda_1=405\text{nm}$  and  $\lambda_2=650\text{nm}$ , it derives from Equation (11) that  $n_e=1.802$ . Consequently, the birefringent material may be chosen where its refractive indices  $n_e$  and  $n_o$  substantially equal 1.802 and 1.5, respectively.

25 In the present description, two refractive indices  $n_a$  and  $n_b$  are substantially equal where  $|n_a - n_b|$  is equal to or less than, preferably, 0.01 and, more preferably, 0.005, where the values 0.01 and 0.005 are a matter of purely arbitrary choice.

In the second embodiment and by way of illustration only the optical record carriers 3, 3' and 3'' are a BD-format disc, a DVD-format disc and a CD-format disc,

respectively. Firstly, the wavelength  $\lambda_1$  is comprised in the range between 365 and 445nm and, preferably, 405nm. The wavelength  $\lambda_2$  is comprised in the range between 620 and 700nm and, preferably, 650nm. The wavelength  $\lambda_3$  is comprised in the range between 740 and 820nm and, preferably, 785nm. Secondly, the numerical aperture  $NA_1$  equals about 0.85 in the reading mode and in the writing mode. The numerical aperture  $NA_2$  equals about 0.6 in the reading mode and is above 0.6, preferably 0.65, in the writing mode. The numerical aperture  $NA_3$  is below 0.5, preferably 0.45. Thirdly, the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are as follows:  $p_1=p_e$ ,  $p_2=p_e$  and  $p_3=p_o$ .

In the second embodiment, the objective lens 17 is a bi-aspherical element.

The objective lens 17 has a thickness of 2.120mm along the Z-axis (direction of its optical axis) and an entrance pupil with a diameter of 4.0mm. The numerical aperture of the objective lens 17 is equal: to 0.85 at the wavelength  $\lambda_1$  (=405nm), to 0.6 at the wavelength  $\lambda_2$  (=650nm), and to 0.45 at the wavelength  $\lambda_3$  (=785nm). The lens body of the objective lens 17 is made of LASFN31 Schott glass with a refractive index equal to 1.9181 at the wavelength  $\lambda_1$  (=405nm), to 1.8748 at the wavelength  $\lambda_2$  (=650nm), and to 1.8664 at the wavelength  $\lambda_3$  (=785nm). The rotational symmetric aspherical shape of the first and second surface of the objective lens 17 are given by the following equation:

$$H(r) = \sum_{i=1}^5 B_{2i} r^{2i} \quad (12)$$

where "H(r)" is the position of the surface along the optical axis of the lens 17 in millimeters,

"r" is the distance to the optical axis in millimeters, and "B<sub>k</sub>" is the coefficient of the k-th power of H(r). The value of the coefficients B<sub>2</sub> until B<sub>14</sub> for the first surface facing the laser are 0.27025467, 0.013621503, 0.0010887228, 0.00025122383, -5.8150037 10<sup>-5</sup>, 2.1911964 10<sup>-5</sup>, -1.965101 10<sup>-6</sup>, respectively. For the second surface facing the optical record carrier the value of the coefficients B<sub>2</sub> until B<sub>14</sub> for the first surface facing the laser are 0.085615362, 0.029034441, -0.031174254, 0.02322335, -0.012032137, 0.0035665564, -0.00044658898, respectively. The free working distance, i.e. the distance between the objective lens 17 and the optical record carrier, is equal: to 1.000mm at the wavelength  $\lambda_1$  (=405nm) for a BD-format disk having a cover layer thickness of 0.1mm, to 0.7961mm at the wavelength  $\lambda_2$  (=650nm) for a DVD-format disk having a cover layer thickness of 0.6mm, and to 0.4446mm at the wavelength  $\lambda_3$  (=785nm) for a CD-format disk having a cover layer thickness of 1.2mm. The cover layer thickness of the disk is made of polycarbonate with a refractive index equal: to 1.6188 at the wavelength  $\lambda_1$  (=405nm), to 1.5806 at the wavelength  $\lambda_2$

(=650nm), and to 1.5731 at the wavelength  $\lambda_3$  (=785nm). It is noted that the objective lens 17 is compatible with the BD-format. In order to make the objective lens suitable for scanning a DVD-format disc and a CD-format disk, spherical aberration arising due to the difference of cover layer thickness and spherochromatism has to be compensated. Spherical aberration can be expressed in the form of the Zernike polynomials. For further information, see e.g. M. Born and E. Wolf, "Principles of Optics," p.469-470 (6<sup>th</sup> ed.) (Pergamon Press) (ISBN 0-08-09482-4). The amount of spherical aberration  $W_{abb}$  arising from the objective lens 17 as designed according to Equation (12) can be determined by ray-tracing as explained above with reference to Fig. 4.

Therefore, in the second embodiment, the stepped profile is further designed for compensating the wavefront aberration  $W_{abb}$  at the wavelengths  $\lambda_2$  and  $\lambda_3$ . Thus the step heights  $h_j$  are to be chosen such that the wavefront modification  $\Delta W_1$  is flat and such that the wavefront modification  $\Delta W_2$  compensates a wavefront aberration  $W_{abb,2}$  for the wavelength  $\lambda_2$  and the wavefront modification  $\Delta W_3$  compensates a wavefront aberration  $W_{abb,3}$  for the wavelength  $\lambda_3$ .

Accordingly, the step heights  $h_j$  are chosen such that both the phase change  $\Delta\Phi_1(p_1)$  is substantially equal zero modulo  $2\pi$  and such that the sums of the wavefront modifications  $\Delta W_2$  and  $\Delta W_3$  and the wavefront aberration  $W_{abb}$  substantially equal zero at the wavelengths at the wavelengths  $\lambda_2$  and  $\lambda_3$ , respectively, where the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  may substantially differ from each other. By way of illustration only, an example of the second embodiment of the stepped profile is described in the following where the stepped profile includes 23 steps.

Firstly, similarly to Table III, Table VI shows the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  introduced by step heights that equal  $qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$  where  $p_1=p_e$  and "q" is an integer. These values are found from Table I where the values  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  are known a step height that equals  $h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$  where  $p_1=p_o$ , i.e. for  $q=1$ .



Table VI:

q	$\Delta\Phi_2(p_2)/2\pi$ $p_2=p_o$	$\Delta\Phi_3(p_3)/2\pi$ $p_3=p_e$
-1	0.377	0.360
0	0.000	0.000
1	0.623	0.640
2	0.246	0.280
3	0.869	0.920
4	0.492	0.560
5	0.115	0.200
6	0.738	0.840
7	0.361	0.480
8	0.984	0.120
9	0.607	0.760

It is noted in Table VI that the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  have 8 and 11 substantially different values modulo  $2\pi$ , respectively. This is consistent with Table II where  $\#\Delta\Phi_2(p_2)=8$  for  $p_2=p_o$  and  $\#\Delta\Phi_3(p_3)=11$  for  $p_3=p_e$ .

It is also noted that, since the polarization  $p_3$  differs from the polarizations  $p_1$  and  $p_2$ , at least three different values of the phases changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  can be chosen, thereby resulting in allowing the design of the stepped profile with a relatively low number of steps, typically less than 40 steps, since a stepped profile with a high number of steps (typically, 50 or more steps) is of less practical use.

Secondly, similarly to Table IV, Table VII shows the "optimized zones" of the step height  $h_j$  ( $=qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$ ) where  $p_1=p_e$  and the values of the phase change  $\Delta\Phi_2(p_2)/2\pi$  and  $\Delta\Phi_3(p_3)/2\pi$ , which are determined from Table III with  $p_2=p_e$  and  $p_3=p_o$  and the wavefront aberration  $W_{abb}$  (see Fig. 4) according to the method known from said article by B.H.W. Hendriks et al.

Table VII also shows, for a step height  $qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$  where  $p_1=p_o$ , the values of the phase change  $\Delta\Phi_2(p_2)$  for approximating the wavefront  $\Delta W_2$  of the type of spherical aberration according to Table VI where  $p_2=p_o$ . Table VII also shows, for a step height  $qh_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$ , the values of the phase change  $\Delta\Phi_3(p_3)$  for approximating the optimized zones according to Table VI where  $p_3=p_e$ . Table VII also shows the corresponding height  $h_j$  (calculated from Equation (4a) where  $p_1=p_o$ ).

Table VII:

	Zones [mm]	q	$h_j(\mu\text{m})$	$\Delta\Phi_2(p_2)$ $p_2=p_0$	$\Delta\Phi_3(p_3)$ $p_3=p_e$
j=1	0.000-0.230	0	0.000	0.000	0.000
j=2	0.230-0.320	5	4.050	0.723	1.257
j=3	0.320-0.400	2	1.620	1.546	1.759
j=4	0.400-0.470	7	5.670	2.268	3.016
j=5	0.470-0.530	4	3.240	3.091	3.519
j=6	0.530-0.580	1	0.810	3.914	4.021
j=7	0.580-0.640	6	4.860	4.637	5.278
j=8	0.640-0.690	3	2.430	5.460	5.781
j=9	0.690-0.750	8	6.480	6.183	7.037
j=10	0.750-0.820	5	4.050	7.006	7.540
j=11	0.820-0.900	2	1.620	7.829	8.042
j=12	0.900-1.150	-1	-0.810	8.652	8.545
j=13	1.150-1.205	2	1.620	7.829	—
j=14	1.205-1.240	5	4.050	7.006	—
j=15	1.240-1.270	8	6.480	6.183	—
j=16	1.270-1.295	3	2.430	5.460	—
j=17	1.295-1.315	6	4.860	4.637	—
j=18	1.315-1.335	1	0.810	3.914	—
j=19	1.335-1.352	4	3.240	3.091	—
j=20	1.352-1.368	7	5.670	2.268	—
j=21	1.368-1.380	2	1.620	1.546	—
j=22	1.380-1.395	5	4.050	0.723	—
j=23	1.395-1.325	3	0.000	-0.823	—

It is noted in Table VII that both the phase changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  associated with the corresponding “optimized zones” approximate a wavefront modification of the type of spherical aberration and defocus. In other words, the optical scanning device provided with the NPS according to Table VII is advantageously compatible with the BD-format, the DVD-format and the CD-format, since it requires only one objective lens.

It is also noted that the polarization  $p_3$  differs from the polarization  $p_1$ , at least three different values of the phases changes  $\Delta\Phi_2(p_2)$  and  $\Delta\Phi_3(p_3)$  can be chosen, thereby resulting in allowing the design of the stepped profile with a relatively low number of steps, typically less than 40 steps, since a stepped profile with a high number of steps (typically, 50 or more steps) is of less practical use.

Fig. 7 shows a curve 83 representing the step height  $h(x)$  of the NPS 24 according to Table VII. It is noted in respect of the curve 83 that the stepped profile is designed such that the relative step heights  $h_{j+1}-h_j$  between adjacent steps include a relative step height having an optical path substantially equal to  $a\lambda_i$ , wherein " $a$ " is an integer and  $a>1$  and " $\lambda_i$ " is the design wavelength. In other words, such a relative step height is higher than the reference height  $h_{ref}(\lambda_{ref}=\lambda_1, p_{ref}=p_1)$ .

Similarly to Table V, Table VIII shows the values  $OPD_{rms}[W_{abb}+\Delta W_i]$  for the wavefront modifications  $\Delta W_1$ ,  $\Delta W_2$  and  $\Delta W_3$  where the radiation beams 15, 15' and 15'' (at the respective wavelengths and polarizations) traverse the NPS according to Table VII (and shown in Fig. 7). Table VIII also shows the values  $OPD_{rms}[W_{abb}]$  associated with the wavefront aberration  $W_{abb}$  (i.e. without the correction of the NPS 24 according to Table VII). The values  $OPD_{rms}[W_{abb}+\Delta W_i]$  and  $OPD_{rms}[W_{abb}]$  have been calculated from ray-tracing simulations.

Table VIII:

	$OPD_{rms}[W_{abb}+\Delta W_i]$	$OPD_{rms}[W_{abb}]$
$i=1 (p_1=p_o)$	$1.1m\lambda$	$1.1m\lambda$
$i=2 (p_2=p)$	$41.3m\lambda$	$466.8m\lambda$
$i=3 (p_3=p_e)$	$64.4m\lambda$	$202.5m\lambda$

It is noted in Table VIII that the three three values  $OPD_{rms}[W_{abb}+\Delta W_i]$  are below the diffraction limit, i.e. less than  $70 m\lambda$ , for the NPS 24 according to Table VII, thereby allowing any format of optical record carriers to be scanned.

As an alternative of the second embodiment of the stepped profile, the value  $\Delta\Phi_2(p_2)$  is substantially equal to the value  $\Delta\Phi_3(p_3)$ , where the polarization  $p_2$  different from the polarization  $p_3$ , i.e.:

$$\Delta\Phi_2(p_2)=\Delta\Phi_3(p_3) \quad (13)$$

In the case where  $p_1=p_o$ ,  $p_2=p_o$  and  $p_3=p_e$  it derives from Equations (0c), (5b), (5c) and (13) that:

$$\frac{\lambda_2}{n_o - 1} = \frac{\lambda_3}{n_e - 1} \quad (14)$$

It follows from Equation (14) that:

$$n_e = 1 + \frac{\lambda_3}{\lambda_2} (n_o - 1) \quad (15)$$

Thus, for example, in the case where  $n_o=1.5$ ,  $\lambda_3=785\text{nm}$  and  $\lambda_2=650\text{nm}$ , it derives from

Equation (15) that  $n_e=1.603$ . Consequently, the birefringent material may be chosen where its refractive indices  $n_e$  and  $n_o$  substantially equal 1.603 and 1.5, respectively.

Whilst in the above described embodiment an optical scanning device compatible with a CD-format disc, a DVD-format disc and a BD-format disc or HD-DVD format disc is described, it is to be appreciated that the scanning device according to the invention can be alternatively used for any other types of optical record carriers to be scanned.

An alternative of the stepped profile described above is designed for introduced a symmetric wavefront modification of a type other than spherical aberration, e.g., of the type of defocus. For more information on the mathematical functions representing such wavefront modifications, see, e.g. the book by M. Born and E. Wolf entitled "Principles of Optics," pp.464-470 (Pergamon Press 6<sup>th</sup> Ed.) (ISBN 0-08-026482-4).

In other alternatives of the stepped profiles described above, the wavelength  $\lambda_2$  or  $\lambda_3$  is chosen as the design wavelength  $\lambda_{\text{ref}}$ . Table IX shows the values of the reference height  $h_{\text{ref}}(\lambda, p)$  in the case where the wavelength  $\lambda_{\text{ref}}$  equals  $\lambda_2$  or  $\lambda_3$  and the polarization  $p_{\text{ref}}$  equals  $p_o$  or  $p_e$  and where, e.g.,  $n_o=1.5$ ,  $n_e=1.62$ ,  $\lambda_2=650\text{nm}$  and  $\lambda_3=785\text{nm}$ .

Table IX:

	$h_{\text{ref}}(\lambda_{\text{ref}}, p_{\text{ref}})$	
	$\lambda_{\text{ref}}=\lambda_2$	$\lambda_{\text{ref}}=\lambda_3$
$p_{\text{ref}}=p_o$	1.300 $\mu\text{m}$	1.570 $\mu\text{m}$
$p_{\text{ref}}=p_e$	1.048 $\mu\text{m}$	1.266 $\mu\text{m}$

An alternative to the NPS arranged on the entrance face of the objective lens may be of any shape like a plane.

As an alternative to the optical scanning device described with wavelengths of 785nm, 660nm and 405nm are used, it is to be appreciated that radiation beams of any other combinations of wavelengths suitable for scanning optical record carriers may be used.

As another alternative to the optical scanning device described with the above values of numerical apertures, it is to be appreciated that radiation beams of any other

combinations of numerical apertures suitable for scanning optical record carriers may be used.

As another alternative of the optical scanning device described above, at least one of the polarizations  $p_1$ ,  $p_2$  and  $p_3$  is switched between a first state and a second state such that the NPS introduces a flat wavefront modification when that polarization is in the first state and a wavefront modification of a type of spherical aberration or defocus when that polarization is in the second state. It is noted that the switching of each of the polarizations  $p_1$ ,  $p_2$  and  $p_3$  is known, e.g., from the European Patent application filed on 07.12.2001 with the application number EP 01204786.6.

Alternatively, at least one of the polarizations  $p_1$ ,  $p_2$  and  $p_3$  is switched between a first state and a second state such that the NPS introduces a first amount of wavefront modification of the type(s) of spherical aberration and/or defocus when that polarization is in the first state and a second, different amount of wavefront modification of the type(s) of spherical aberration and/or defocus when that polarization is in the second state.

In a particular case, each of the polarizations  $p_1$ ,  $p_2$  and  $p_3$  is switched between a first state and a second state such that the NPS introduces a flat wavefront modification when the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the first states and a wavefront modification of the type(s) of spherical aberration and/or defocus when the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the second states. This advantageously allows to design the NPS for introducing, in respect of the wavelengths  $\lambda_1$ ,  $\lambda_2$  and  $\lambda_3$ : three respective flat wavefront modifications when the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the first states, respectively, and three wavefront modifications of the type(s) of spherical aberration and/or defocus when the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the second states, respectively. Accordingly, the NPS has no optical effect where the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the first states and has an optical effect (by generating wavefront modifications of the type(s) of spherical aberration and/or defocus) where the polarizations  $p_1$ ,  $p_2$  and  $p_3$  are in the second states.

It is noted in respect of the above that the polarizations  $p_1$ ,  $p_2$  and  $p_3$  can be switched independently so that the optical scanning device provided with such a NPS has eight different configurations.